

## Synthesis and Structure of Ta<sub>4</sub>SI<sub>11</sub>: Disorder and Mixed Valency in the First Tantalum Sulfide Iodide

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The new compound Ta<sub>4</sub>SI<sub>11</sub> has been prepared by direct reaction of the elements at 430 °C for 2 weeks in evacuated Pyrex ampules and characterized by single-crystal X-ray diffraction, X-ray photoelectron spectroscopy, magnetic susceptibility measurements, and semiempirical electronic structure calculations. Ta<sub>4</sub>SI<sub>11</sub> crystallizes with orthorhombic symmetry in space group *Pmmn*; *a* = 16.135(3) Å, *b* = 3.813(1) Å, *c* = 8.131(2) Å, and *Z* = 1. The disordered structure involves two crystallographically distinct sites for Ta atoms, both of which are 50% occupied as well as a bridging anion site that is 50% S and 50% I. Magnetic susceptibility above 100 K gives  $\mu_{\text{eff}} = 1.53 \mu_{\text{B}}$  to suggest one unpaired electron per formula unit. X-ray photoelectron spectroscopy and extended Hückel calculations suggest that the structure consists of Ta<sub>3</sub> triangles and "isolated" Ta atoms, leading to the formulation (Ta<sub>3</sub>)<sup>9+</sup>(Ta<sup>4+</sup>)(S<sup>2-</sup>)(I<sup>-</sup>)<sub>11</sub> and we hypothesize that each Ta<sub>3</sub> is capped by a sulfur atom.

### Introduction

The number of well-characterized ternary tantalum chalcogenide halides is meager, limited to pseudo one-dimensional chalcogen-rich compounds such as TaTe<sub>4</sub>I,<sup>1</sup> (TaSe<sub>4</sub>)<sub>n</sub>I (*n* = 2, 3),<sup>2,3</sup> and Ta<sub>4</sub>Se<sub>16</sub>Br<sub>n</sub> (*n* = 1, 2),<sup>3,4</sup> and to the layered halide-rich compounds Ta<sub>3</sub>SBr<sub>7</sub><sup>5</sup> and Ta<sub>3</sub>QI<sub>7</sub> (Q = Se, Te),<sup>6</sup> recently discovered in our laboratory. A few others have been mentioned in the literature, with only cursory structural characterization. TaS<sub>2</sub>Cl<sub>2</sub>, TaS<sub>2</sub>Cl<sub>2</sub>, TaS<sub>2</sub>Cl<sub>3</sub>, TaSBr<sub>3</sub>, and TaSeBr<sub>3</sub><sup>7</sup> have been assigned their various structural features on the basis of chemical analysis and metal-chalcogen stretching frequencies only. No tantalum sulfide iodide, to the best of our knowledge, has ever been reported. Certainly, the tally of ternary tantalum chalcogenide halide compounds lags well behind that of its lighter relative niobium, which has yielded a rich ternary chemistry. NbTe<sub>4</sub>I,<sup>8</sup> NbSeX<sub>3</sub>,<sup>9,10</sup>

Nb<sub>3</sub>QX<sub>7</sub>,<sup>11–16</sup> NbSeI,<sup>17,18</sup> and Nb<sub>6</sub>SI<sub>9</sub><sup>19</sup> display remarkable structural diversity, adopting structures with isolated niobium atoms, Nb–Nb dimers, Nb<sub>3</sub> triangular clusters, Nb<sub>4</sub> tetrahedral clusters, and Nb<sub>6</sub> octahedral clusters, respectively. Additionally, centered Nb<sub>6</sub> trigonal prismatic clusters in the case of the alkali-metal salts A<sub>3</sub>Nb<sub>6</sub>SBr<sub>17</sub>,<sup>20,21</sup> and tetranuclear butterfly clusters in CsNb<sub>4</sub>Cl<sub>11</sub>,<sup>22</sup> have been characterized. Generally, the metal atoms in these compounds can be assigned a single oxidation state (though occasionally a

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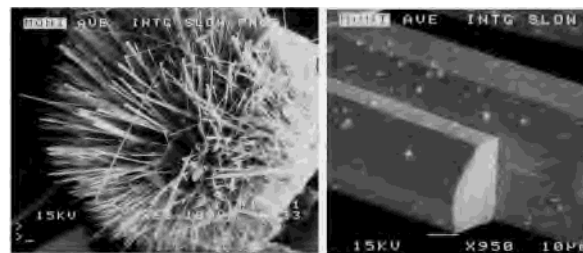
nonintegral one). Compounds made at higher temperatures seldom exhibit mixed-valent behavior, where the same metal exists in two different formal oxidation states in the same compound. A few examples from niobium chemistry are known, however. The binary halide  $\text{Nb}_3\text{Cl}_8$  actually forms in the composition range  $\text{Nb}_{3-x}\text{Cl}_8$  ( $x \leq 0.44$ )<sup>23</sup> and has been described as a mixed  $(\text{Nb}^{4+})_2\text{Cl}_8/(\text{Nb}^{2.67+})_3\text{Cl}_8$  system.<sup>12</sup>  $\text{Nb}_3\text{Se}_5\text{Cl}_7$  has been formulated  $(\text{Nb}^{5+})(\text{Nb}^{4+})_2(\text{Se}_2^{2-})_2(\text{Se}^{2-})(\text{Cl}^-)_7$ .<sup>24</sup> Another example is the peculiar compound  $\text{Nb}_7\text{S}_2\text{I}_9$ , where trigonal bipyramidal  $(\text{Nb}^{5+})_5$  molecules are encapsulated inside a layered  $(\text{Nb}^{3+})_3\text{SI}_7$  framework.<sup>25</sup> In this compound, however, the  $\text{NbI}_5$  units can almost be considered a kind of molecule of crystallization since they are encapsulated in large channels and are not directly joined to the extended  $\text{Nb}_3\text{SI}_7$  framework. Such compounds, besides providing structural novelty, are also interesting examples of the delicate balance possible between competing redox forces within a structure.<sup>26</sup>

During the course of our studies of the ternary chalcogenide halide chemistry of tantalum, we have focused on halide-rich compositions, usually near  $\text{Ta}_3\text{QX}_7$ . As part of an attempt to synthesize  $\text{Ta}_3\text{SI}_7$ , we have discovered the novel mixed-valent compound  $\text{Ta}_4\text{SI}_{11}$ , containing both trinuclear  $\text{Ta}_3^{3+}$  clusters and isolated  $\text{Ta}^{4+}$  ions amid undulating, disordered sulfur and iodine layers. Herein, we report the synthesis as well as structural and theoretical characterization of this new structure type.

## Experimental Section

**Synthesis.**  $\text{Ta}_4\text{SI}_{11}$  was first observed as an abundant (ca. 50%) side product from reactions of the elements in the molar ratio 3:1:7 Ta:S:I at 450 °C in sealed, evacuated borosilicate glass tubes, designed to grow  $\text{Ta}_3\text{SI}_7$  phases. Because of the inability to identify the compound from its Guinier powder diffraction pattern, a single crystal was selected and its structure solved (see below), yielding the composition “ $\text{Ta}_4\text{SI}_{11}$ ”. Tubes were then loaded at this stoichiometry and heated at various temperatures to determine optimum synthetic conditions.  $\text{Ta}_4\text{SI}_{11}$  is best made by stoichiometric reaction of the elements at 430 °C, for a duration of 2 weeks. Growth of the compound is sensitive to temperature gradients in the tube: highest (nearly quantitative) yields were obtained when the tube furnace was packed tightly with asbestos, to smooth out uneven temperature regions as much as possible. Heating a  $\text{Ta}_4\text{SI}_{11}$  mixture with an imposed temperature gradient leads to a mixture of phases, including hexagonal and orthorhombic  $\text{Ta}_3\text{SI}_7$ , and  $\text{Ta}_6\text{I}_{14}$  in addition to  $\text{Ta}_4\text{SI}_{11}$ . The reagents used were tantalum “turnings” (Aesar, 99.99%, cleaned with an  $\text{HF}/\text{HNO}_3/\text{H}_2\text{SO}_4$  solution to remove surface impurities, then rinsed with ethanol, and dried in vacuo at 500 °C), sulfur powder (Alfa, 99.9%, sublimed twice before use), and iodine (Alfa, 99.9%, resublimed). All purified reagents and products were handled in an Ar-filled glovebox.

$\text{Ta}_4\text{SI}_{11}$  crystallizes as long silver bars, which often aggregate in dense thickets (Figure 1). The brittle crystals splinter easily along the long axis.  $\text{Ta}_4\text{SI}_{11}$  is sensitive to oxygen and moisture,



**Figure 1.** SEM images of (left) a typical spiny ball formation of  $\text{Ta}_4\text{SI}_{11}$ ; (right) higher magnification of the end of a crystal.

decomposing to an uncharacterized white powder after several days. In inert atmospheres, the compound appears to be indefinitely stable toward decomposition up to temperatures of  $\approx 525$  °C, above which only  $\text{Ta}_6\text{I}_{14}$ ,  $\text{TaI}_5$ , and  $\text{TaS}_2$  are observed by X-ray powder diffraction.

**Scanning Electron Microscopy.** Verification of the presence of tantalum, sulfur, and iodine in multiple single-crystal samples of  $\text{Ta}_4\text{SI}_{11}$  was carried out using a JEOL 6100 scanning electron microscope operating at 15 kV.

**X-ray Photoelectron Spectroscopy.** XPS spectra were measured using a PHI 5500 multitechnique surface analysis equipment instrument, using Mg  $K\alpha$  radiation (29.35 eV). Samples were prepared by finely powdering several larger single crystals of  $\text{Ta}_4\text{SI}_{11}$  in the glovebox immediately before use. The samples were pressed onto indium foil in the glovebox and transferred to the XPS instrument via an inert atmosphere carrier. Binding energies were corrected for charging using the C 1s peak of adventitious carbon (284.8 eV) as an internal reference.<sup>27</sup> However, no significant charging effects were observed, as C 1s always appeared at 284.8 eV. Additionally, large intact single crystals were mounted and etched by sputtered Ar ions to remove any surface impurities introduced during handling, to probe the presumably pristine interior of the sample.

**Magnetic Susceptibility.** Temperature-dependent magnetic susceptibility measurements were performed with a Quantum Design SQUID magnetometer. Aggregations of  $\text{Ta}_4\text{SI}_{11}$  bars were loaded and sealed under an inert atmosphere in fused silica tubes. The samples were kept in place by means of two tightly fitting fused silica rods on either side of the sample. The samples were chosen so that only larger single-crystal aggregations were used, to avoid unwanted powder impurities and to minimize the surface area susceptible to oxidation. Measurements were taken from 4 to 300 K, at a field strength of 3 T. The results are shown in Figure 2. A room-temperature moment of  $1.53 \mu_B$  was measured, reproducible over three separate runs.

**X-ray Crystallography.** A small silver bar ( $0.02 \times 0.03 \times 0.12 \text{ mm}^3$ ) was sealed inside a glass capillary under argon and mounted on a Siemens P4 diffractometer. A primitive orthorhombic unit cell with dimensions  $a = 16.135(3) \text{ \AA}$ ,  $b = 3.813(1) \text{ \AA}$ , and  $c = 8.131(2) \text{ \AA}$  was indexed and refined on the basis of 45 reflections in the range  $6^\circ \leq 2\theta \leq 25^\circ$ . Axial photographs confirmed these lattice metrics. 2629 reflections were collected to  $2\theta_{\text{max}} = 55^\circ$ , of which 685 were unique, 533 unique observed ( $I > 2\sigma I$ ), and  $R_{\text{int}} = 0.0498$ . An empirical absorption correction was applied to the data, using the average of six complete “psi-scans” measured in well-separated regions of reciprocal space. The structure was solved using direct methods (SHELX-86) and refined with the SHELXL-93 crystal-

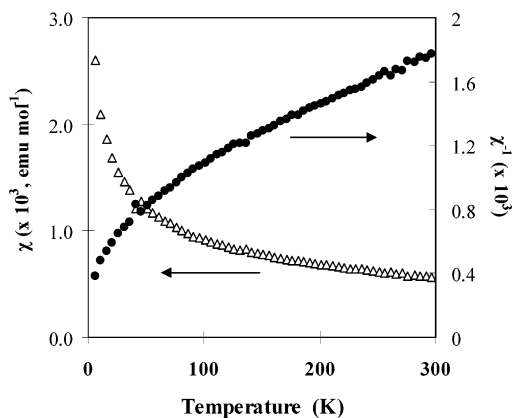
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**Figure 2.** Temperature-dependent magnetic susceptibility for  $Ta_4SI_{11}$  (molar susceptibility  $\chi$ ,  $\text{emu mol}^{-1}$ , and inverse molar susceptibility  $\chi^{-1}$ ,  $\text{mol emu}^{-1}$ ). A fit to Curie-Weiss behavior for temperatures exceeding 100 K gives  $\mu_{\text{eff}} = 1.53 \mu_B$ .

**Table 1.** Summary of Crystallographic Data for  $Ta_4SI_{11}$

fw	2151.76
space group	$Pmmn$ (No. 59)
lattice params	
$a$ (Å)	16.135(3)
$b$ (Å)	3.813(1)
$c$ (Å)	8.131(2)
vol (Å <sup>3</sup> )	500.2(2)
$Z$	1
$d_{\text{calc}}$ ( $\text{g cm}^{-3}$ )	7.143
radiation	Mo $K\alpha$ , $\lambda = 0.71071$ Å
linear abs coeff	38.854 $\text{mm}^{-1}$
temp	23 °C
residuals <sup>a</sup>	
R1 (observed/all data)	0.0322/0.0503
wR2 (observed/all data)	0.0695/0.0762

$$^a R1 = \frac{\sum ||F_o| - |F_c||}{\sum |F_o|}; wR2 = \frac{[\sum w(F_o^2 - F_c^2)^2 / \sum w(F_o^2)]^{1/2}}{w}; w = 1/\sigma^2(F_o).$$

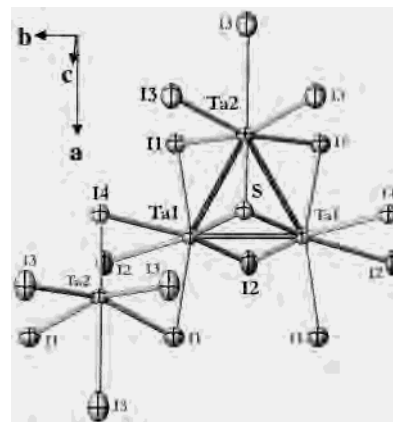
**Table 2.** Important Internuclear Distances (Å) in  $Ta_4SI_{11}$

atom-atom	distance (Å)
Ta1-Ta1	2.998(2)
Ta1-Ta2	3.157(1)
Ta1-I1	2.769(1) × 2
Ta1-I2	2.583(2); 3.127(2)
Ta1-I4	2.892(2)
Ta1-S	2.523(2)
Ta2-I1	2.744(1) × 2
Ta2-I3	2.780(1) × 2; 3.112(2)
Ta2-I4	2.286(1)
Ta2-S	2.324(5)

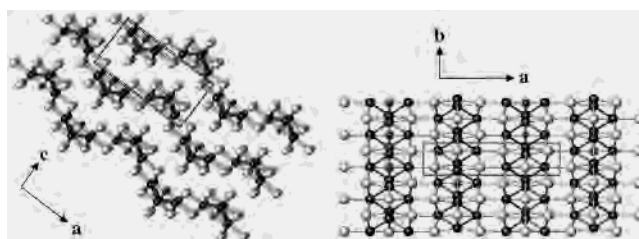
lographic package.<sup>28,29</sup> Further data collection details and important internuclear distances are given in Tables 1 and 2, respectively.

## Results and Discussion

Figure 3 shows a 80% probability displacement ellipsoid plot of the local Ta environments in  $Ta_4SI_{11}$ , representing the  $(Ta_3)^{9+}/Ta^{4+}$  disorder model discussed below. Figure 4a shows  $Ta_4SI_{11}$  viewed parallel to the short  $b$  axis. The compound can be viewed as a close-packed structure similar to the  $M_3QX_7$  compounds, but with a different metal atom distribution. In  $Ta_4SI_{11}$ , approximately close-packed layers



**Figure 3.** Displacement ellipsoid view (80% probability) and atom labeling scheme showing the  $Ta_3$  clusters and isolated Ta centers hypothesized to be disordered throughout the structure of  $Ta_4SI_{11}$ . Independent atomic positions are labeled in bold; symmetry-equivalent positions are smaller and gray.



**Figure 4.** (a) [010] projection of  $Ta_4SI_{11}$ . Larger black circles are Ta atoms; smaller black circles are S atoms; large white circles are I atoms. (b) [001] projection of  $Ta_4SI_{11}$ . Both figures illustrate all fully and partially occupied positions. In the [001] projection, it is clear that two Ta sites cannot be simultaneously occupied.

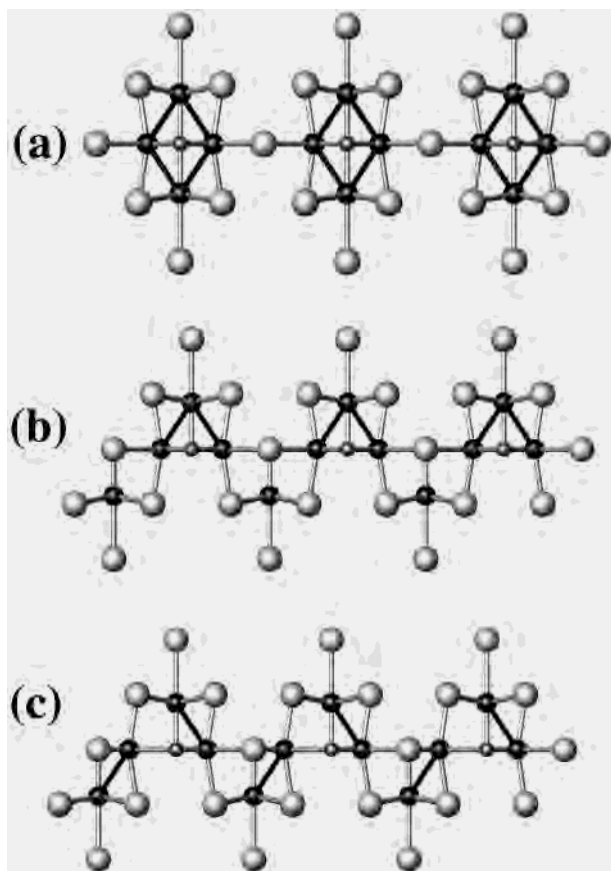
of composition “ $SI_{11}$ ” stack parallel to the  $(\bar{3}02)$  family of lattice planes in cubic close-packed ( $\dots ABC\dots$ ) or ( $\dots C\dots$ ) fashion, at an angle of  $37.25^\circ$  from the  $a$  axis.

However, instead of a pattern of interstitial site occupation by the metal atoms that generates flat parallel slabs as in the layered  $M_3QX_7$  compounds, in  $Ta_4SI_{11}$  the occupation of these sites defines an undulating sheet structure with a period of 16.135 Å.  $Ta_4SI_{11}$  can still be viewed as a “layered” structure since the undulating  $Ta_4SI_{11}$  layers are separated by a van der Waals gap.

The nature of the distribution of tantalum atoms within the layers is not straightforward. The structure solution gives rise to two crystallographically distinct tantalum positions. Full occupation of these sites would lead to butterfly  $Ta_4$  clusters that bend at the turns of the undulating anion layers. However, the site occupation factors of both metal sites each refine to near 50%. An additional complication is the atomic position at the cradle or elbow of the apparent bent cluster, seemingly bridging all four metals. This site refines as a 50:50 sulfur-iodine split position (S/I4). Figure 4b shows a [001] view of what the crystallographic experiment sees, which is an average over the entire crystal. The question now becomes how the metal atoms are distributed throughout the structure and how the metal atom distribution correlates with the split S/I4 position. We can attribute the various site occupancies to structurally ordered ribbons of Ta clusters, shown in Figure 5, with the distribution of ribbons along [100] as well as along different layers in the layer stacking

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**Figure 5.** Three possible distributions of Ta atoms that may exist in  $\text{Ta}_4\text{SI}_{11}$ . (a) “Butterfly” clusters; (b) triangles alternating with isolated Ta atoms; (c) herringbone pattern of  $\text{Ta}_2$  dimers.

(001) direction being completely disordered. Within a ribbon along [010], there are three possibilities for the distribution of Ta atoms (see Figure 5): (1)  $\text{Ta}_4$  “butterfly” clusters linked by bridging S/I4 ions; (2) alternating  $\text{Ta}_3$  triangular clusters and isolated Ta atoms; and (3) a “herringbone” pattern of  $\text{Ta}_2$  dimers. The distribution of S and I4 atoms should be closely correlated with the patterns of Ta clusters in possibilities (1) and (2), but less so for possibility (3). The solution to this problem cannot be conclusively determined based on the X-ray diffraction experiment, which measures average electron density over the entire crystal and cannot elucidate the precise local situation: independent characterization methods are required.

**X-ray Photoelectron Spectroscopy.** To help clarify more details about the structure, X-ray photoelectron spectroscopy spectra were measured. The XPS spectrum of any tantalum atom produces a doublet signal, corresponding to the  $4f_{7/2}$  (lower binding energy) and  $4f_{5/2}$  (higher binding energy). These peaks are always separated by 1.91 eV. This constant separation facilitates deconvolution of overlapping peaks due to inequivalent atoms. The binding energies reported below are for the Ta  $4f_{7/2}$  peaks. The  $\text{Ta}_4\text{SI}_{11}$  spectrum showed two peaks at 23.8 and 26.3 eV with visibly different integrated intensities (the intensity of the 23.8-eV peak is noticeably greater than that of the 26.3-eV peak), which suggest two different oxidation states for Ta in different concentrations. The Ta  $4f_{7/2}$  binding energy of 23.8 eV indicates a reduced

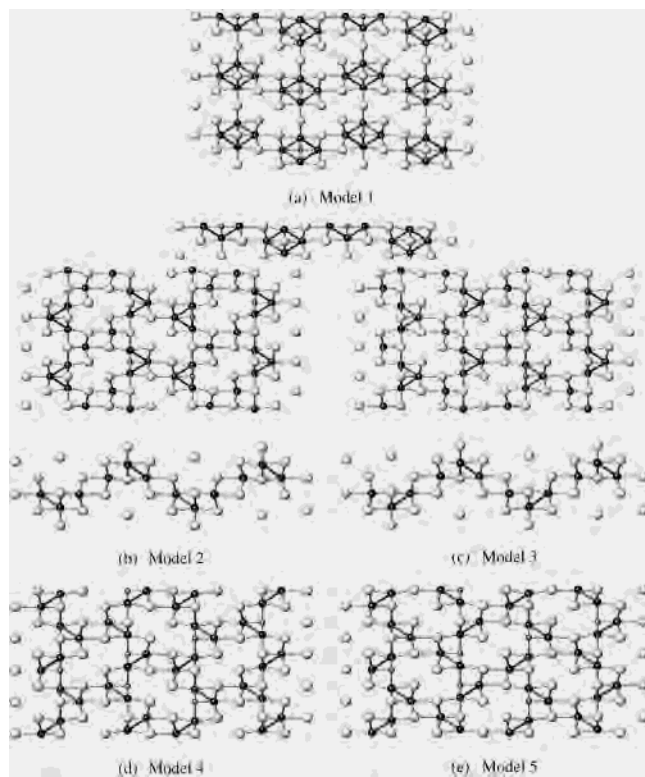
**Table 3.** Tantalum  $4f_{7/2}$  Binding Energies (eV) Measured by X-ray Photoelectron Spectroscopy

compound	binding energy (eV)
Ta metal	21.6–21.9
$(\text{NEt}_4)_2[\text{Ta}_6\text{Cl}_{12}]\text{Cl}_6$	23.8
$[\text{Ta}_6\text{Cl}_{12}]\text{Cl}_2(\text{H}_2\text{O})_4 \cdot 4\text{H}_2\text{O}$	25.8
$\text{K}_3\text{TaO}_4$	25.9
$\text{TaS}_2$	26.7
$\text{Ta}_2\text{O}_5$	26.7
$\text{TaBr}_5$	26.9
$\text{K}_2\text{TaF}_7$	29.4
<i>hc</i> - $\text{Ta}_3\text{SeI}_7$ <sup>6</sup>	23.4
<i>hc</i> - $\text{Ta}_3\text{TeI}_7$ <sup>6</sup>	23.4

metal center and suggests the presence of triangular clusters having an average oxidation state of +3. The binding energies of Ta in *hc*- $\text{Ta}_3\text{SeI}_7$  and *hc*- $\text{Ta}_3\text{TeI}_7$ , which contain  $\text{Ta}_3$  clusters, are 23.4 eV.<sup>6</sup> For comparison, literature XPS values<sup>27</sup> for tantalum in other compounds are listed in Table 3. It is also possible to assign the peak at 26.3 eV to oxygen contamination of the sample, and not to a mixed-valence effect within the compound. However, oxygen contamination of the crystalline samples used would be limited to surface oxidation since all products were handled under an inert atmosphere. To explore this possibility, samples of single crystals were etched in the XPS instrument by sputtering with Ar ions. Etching removes contaminated surface layers and exposes the nonoxidized, presumably pristine interior of the sample. After etching, the peak at 26.3 eV persisted, which suggests this peak is characteristic of the sample, and not due to surface oxidation. Additionally, the Ta  $4f_{7/2}$  peak of  $\text{Ta}_2\text{O}_5$  appears at slightly higher binding energy (26.7 eV, Table 3).

The two binding energies suggest a mixture of  $\text{Ta}_3$  trinuclear clusters and tantalum atoms in a higher oxidation state. It is not likely that the two peaks result from chemical inequivalencies within the trinuclear cluster itself (local symmetry  $C_3$ ), as the 2.8-eV difference between the two is much too large to support this contention. The distortion of the  $\text{Ta}_3$  clusters away from perfect equilateral geometry is moderate, with Ta–Ta distances of 3.157(2) and 2.999(2) Å and angles of 56.70(1)° and 61.65(1)°. A formulation of  $\text{Ta}_4\text{SI}_{11}$  consistent with the crystallographic, magnetic, and X-ray photoelectron spectroscopic data is “ $(\text{Ta}^{3+})_3(\text{Ta}^{4+})(\text{S}^{2-})(\text{I}^-)_{11}$ ”; a 1:1 mixture of  $\text{Ta}_3$  clusters and isolated Ta atoms in the +4 oxidation state. These two metal species must be randomly distributed throughout the three-dimensional structure to give the apparent averaged structure solved from the X-ray diffraction data. While it is possible (but, as yet, unverifiable) that some kind of ordering of the trinuclear clusters and isolated atoms exists within each undulating layer, certainly no ordering or registry can exist between the  $\text{Ta}_4\text{SI}_{11}$  layers; otherwise, different symmetry, unit cell, or superstructure reflections would be found. This formulation is also bolstered by the magnetic data, where the derived magnetic moment of 1.53  $\mu_B$  indicates one unpaired electron per formula unit, close to the spin-only value for isolated  $d^1 \text{Ta}^{4+}$  centers.

The assignment of the split sulfur–iodine position remains, but we hypothesize an arrangement commensurate with the pattern of trinuclear clusters given our experience with



**Figure 6.** Five different structural models of Ta atom distribution used for extended Hückel calculations of total energies.

various  $M_3YX_7$  compounds. Therefore, at each trinuclear cluster, sulfur is the capping atom. Conversely, then, when the single tantalum atom occurs, it is surrounded by an octahedral coordination environment of entirely iodine. This assignment accounts for the amount of each species present in this structure, as well as being the most reasonable placement of the sulfur atoms in light of the behavior of the chalcogen in the similar compounds  $Ta_3QX_7$ ,<sup>6</sup> that is, always capping a trinuclear cluster. While most Ta–I and Ta–S bond distances are typical,<sup>5,6</sup> there is, however, an extremely short Ta2–I4 bond length of 2.286(1) Å, around the  $Ta^{4+}$  ions. However, with the disorder and fractional occupancy present throughout the structure, such a distance could be a false artifact forced upon the structure as a consequence of averaging. No unusual motion is manifest in the anisotropic displacement parameters for S or I4 (see Supporting Information). The only noticeably large  $U_{ij}$  value is  $U_{33}$  for I2, apparently elongated into the van der Waals gap.

**Extended Hückel Calculations.** Extended Hückel calculations were carried out on five structural models (illustrated in Figure 6a–e) of the  $Ta_4SI_{11}$  structure based upon the three models enumerated above, using atomic orbital parameters given in Table 4. The purpose of these calculations was to determine the total energies of various possible metal atom distributions. The unit cells for all five models utilized doubled  $b$  axis lengths to account for the Ta positions along the ribbons but kept the observed values of  $a$  and  $c$  axes lengths. The five models are (view normal to the layers in all cases with additional [010] views for 6b and 6c):

Model 1: “Butterfly”  $Ta_4$  clusters, distributed over the three-dimensional structure as shown in Figure 6a. Along

**Table 4.** Atomic Orbital Parameters Used in Extended Hückel Calculations on Different Models of  $Ta_4SI_{11}$

atom	orbital	$H_{ii}$ (eV)	$\zeta_1$	$c_1$	$\zeta_2$	$c_2$
Ta	6s	−8.96	2.28			
	6p	−4.99	2.24			
	5d	−9.83	4.76	0.6104	1.94	0.6104
S	3s	−20.00	2.12			
	3p	−13.30	1.83			
I	5s	−18.00	2.66			
	5p	−12.72	2.32			

the [010] direction of each undulating  $Ta_4SI_{11}$  layer, a  $Ta_4$  cluster is fully occupied, with the adjacent  $Ta_4$  set unoccupied. For this model, the sulfur was positioned in the crook of the cluster, coordinated to all four Ta atoms.

Models 2 and 3: A 50:50 mixture of  $Ta_3$  clusters and isolated Ta atoms. There are two possible unique distributions in this case: (Figure 6b) such that the  $Ta_3$  clusters are all parallel throughout the structure when viewed along [010], or (Figure 6c), the  $Ta_3$  clusters alternate orientation, following the undulating layers.

Models 4 and 5: A “herringbone” pattern of  $Ta_2$  dimers along [010], with each dimer bridged by a sulfur atom and an iodine atom (Figure 6d,e).

The calculations showed a significant energetic separation among the three broad cases: the total energies of the two  $Ta_3 + Ta$  cases (Figure 6b,c) were found to be quite similar, differing by only 0.07 eV (per formula unit) from one another. However, both of these cases were clearly stabilized relative to the butterfly clusters by 0.27 eV per formula unit and to the dimers by 1.57 eV.

$Ta_4SI_{11}$  bears a structural resemblance to the novel niobium sulfide iodide  $o-Nb_3SI_7$ .<sup>11</sup> Both structures contain undulating mixed sulfur–iodine layers, with the metal atoms defining a similar undulating layer structure. Within both undulating layer structures, triangular clusters are capped by a sulfur atom, generating local  $M_3X_{13}$ -type clusters. In  $o-Nb_3SI_7$ , the  $Nb_3$  cluster coordination environment involves an additional sulfur atom; this second coordinating S provides a bridge to another  $Nb_3$  cluster. Coordination of the  $Nb_3$  units is then “ $Nb_3S_2I_{11}$ ”. In  $Ta_4SI_{11}$ , the sulfur atom is not shared between separate clusters, and the  $Ta_3$  coordination is “ $Ta_3SI_{12}$ ”. Because of these structural similarities, and considering the crystallographic difficulties encountered in the solutions of both structures, one might raise the possibility of whether  $Ta_4SI_{11}$  might actually be “ $o-Ta_3SI_7$ ”, isostructural with the Nb case, but with a disorder mechanism falsely but reproducibly generating two different solutions. In fact, we have also synthesized a compound whose powder pattern is virtually the same as that of  $o-Nb_3SI_7$ , though diffraction quality single crystals have not been grown.<sup>30</sup> Examination of Guinier X-ray powder diffraction patterns of pure samples of  $Ta_4SI_{11}$  and “ $o-Ta_3SI_7$ ” as well as the powder pattern from the published  $o-Nb_3SI_7$  structure was definitive in establishing the uniqueness of the two structure types. It is clear that the patterns are significantly different, reflecting the different structures that emerge from the single-crystal structure determinations.

(30) Smith, M. D.; Miller, G. J. Unpublished results.

The available data support the contention of a mixed-valent, disordered sulfide iodide. The problem of the short Ta<sub>2</sub>–I<sub>4</sub> distance remains and is a disturbing one. Other attempts at characterizing the sample have been attempted: Various mass spectrometry techniques (electron impact ionization and chemical ionization in positive and negative ion modes) have been employed, but the compound is not sufficiently volatile to produce any vapor-phase fragments that could illuminate structural details. The compound is not soluble without decomposition, and therefore other ionization techniques (electrospray, etc.) cannot be used. The potential of an undetected superstructure exists. In particular, the split

position might be “frozen” or resolved into one position, giving rise to an ordered structure. Examination of several crystals at low temperature, using both a rotating anode X-ray source and CCD area detector data, did not reveal either a different unit cell or superstructure reflections that would indicate a structural change from that reported herein.

**Supporting Information Available:** X-ray crystallographic data in CIF format. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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